

# Investigating the NCQ scaling of elliptic flow at LHC with AMPT

#### Speaker: Liang Zheng Central China Normal University

Eur. Phys. J. A53, 124, (2017) LZ, Hui Li, Hong Qin, Qi-Ye Shou and Zhong-Bao Yin

### Outline

- Elliptic flow and quark number scaling
- Quark coalescence model
- Insights to v<sub>2</sub> NCQ scaling from AMPT
- Outlook and Summary

### Elliptic flow in heavy ion collisions

- Elliptic flow: Momentum space anisotropy of particle production  $v_2 = \langle \cos(2(\phi - \Psi)) \rangle$
- Arising due to coordinate space anisotropy (b>0) and interactions among medium constituents
- Why do we care:
  - Sensitive to the properties of QGP medium, e.g. shear viscosity
  - Provide information on the initial state of the collisions



## v<sub>2</sub> and quark number scaling



- Mass ordering at low  $p_T$  driven by the hydrodynamic pressure gradient.
- Baryon meson ordering in the high  $p_T$  region.
- NCQ scaling observed in a wide range of KE<sub>T</sub> indicates the dominance of partonic degrees of freedom.

### Quark coalescence model



Assume quark momentum distribution:

$$f_i(\vec{p_{\perp}}) \equiv rac{dN}{d^3p} = g_i(p_{\perp}) \left[1 + 2v_{2,i}(p_{\perp})\cos(2\phi)
ight]$$

Yield simple relations of quark and hadron flow:

$$v_{2,M}(p_{\perp}) \approx 2v_{2,q}(\frac{p_{\perp}}{2}) \qquad v_{2,B}(p_{\perp}) \approx 3v_{2,q}(\frac{p_{\perp}}{3})$$



Workshop on AMPT for Heavy Ion Collisions

### Violation to NCQ scaling



- Deviations from the NCQ scaling at the level of 20% observed in LHC data
- Possible reasons:
  - Narrow wave function width, Resonance decay, Space-Momentum correlations at freeze-out, Phase-space density

### The paradigm in AMPT



### Identified particle v<sub>2</sub> in AMPT



 $v_2$  obtained with event plane method.

Ampt-v1.26t5-v2.26t5 a=0.3, b=0.15 GeV<sup>-2</sup>  $\mu$  =2.3 fm<sup>-1</sup>,  $\alpha_s$ =0.33,  $\sigma_{parton}$ =3 mb PRC 90, 014904 (2014)

$$f(z) \propto z^{-1}(1-z)^a \exp(-b \ m_{\perp}^2/z)$$
$$\sigma_{PP} \approx \frac{9\pi\alpha_s^2}{2\mu^2}$$

Identified particle  $v_2$  within  $|\eta| < 0.8$ agree with data except in the high  $p_T$ region, where model results are systematically smaller than data.

# Charged particle v<sub>2</sub> with different initial and final state conditions



	а	b (GeV <sup>-2</sup> )
Set A	0.3	0.15
Set B	0.5	0.9

 $\kappa \propto 1/[b(2+a)]$  $\kappa_B \approx 1/6\kappa_A$ 

- Small string tension leads to large low p<sub>T</sub> v<sub>2</sub>
- Large parton-parton scattering cross section generates stronger high p<sub>T</sub> v<sub>2</sub>.

## v<sub>2</sub> with different initial conditions



Right after coalescence

- v<sub>2</sub> from different initial conditions converge at high p<sub>T</sub>.
- Mass ordering exists in both two sets, while the magnitude of the mass splitting changes a little.

### Test of scaling properties

PbPb 2.76 TeV 30-40%, hadron formed right after coalescence



- NCQ scaling behavior relies on the initial conditions even within the quark coalescence context.
- Violation to the NCQ scaling observed in the case generated with smaller string tension.
- The final state parton scattering effect is less important in the formation of NCQ scaling phenomenon.

#### Parton distribution before coalescence

Quark distribution before hadronization in PbPb 2.76 TeV 30-40%



- Quark number densities are quite similar in all cases.
- Smaller string tension generates softer initial parton spectrum.
- The violation to NCQ scaling may arise from the higher parton-parton interaction rate in the overlap region.

# v<sub>2</sub> NCQ scaling varying with collision centrality



- NCQ scaling violated in central collisions and restored in peripheral collisions.
- Similar trend can be extracted from the measured data.

# Integrated v<sub>2</sub> varying with collision centrality

PbPb 2.76 TeV,  $v_2$  integrated over 0.3< $(m_T-m_0)/n_a$ <1 GeV



- Integrated flow reaches maximum in semi-central collisions.
- K over  $\pi$  scaled v<sub>2</sub> ratio slightly depends on the centrality.
- p over  $\pi$  scaled v<sub>2</sub> ratio grows from central to peripheral collisions.

#### Hadronic evolution effects

PbPb 2.76 TeV 30-40%



- Primordial v2: formed right after coalescence procedure.
- Resonance decay and hadron rescattering modifies the primordial v<sub>2</sub> in opposite way.
- Responses to the modification of hadronic evolution depend on the particle types.

#### Impact of hadronic evolution to NCQ scaling



Workshop on AMPT for Heavy Ion Collisions

### Energy dependence of NCQ scaling

PbPb 5.02 TeV



 Similar centrality and hadronic evolution impact on NCQ scaling expected with higher collision energy.

#### Outlooks – NCQ scaling in higher order flow



- NCQ scaling seems to work better for v<sub>3</sub>
- Additional constraint on medium expansion

# Outlooks – NCQ scaling in small systems



• NCQ scaling holds better in pPb than in peripheral PbPb

#### Summary

- It is shown in the AMPT framework that NCQ scaling structure not only depends on the hadronization procedure but also relies on the parton dynamics at the initial stage.
- A sizable distortion to NCQ scaling arises due to the hadronic interactions.
- The coalescence AMPT coalescence in coordinate space, what if in momentum space?

#### Thank you for your attention!



#### Why do we care about flow?



- Sensitive to the properties of QGP, e.g.: shear viscosity over entropy density,  $\eta/s,$  of the produced medium
- Provide information on the initial state of the collisions

#### Quark coalescence model

EPJC 62, 237 (2009)



Freeze-out hypersurface  

$$E\frac{dN_{M}(\vec{p})}{d^{3}p} = \int d\sigma^{\mu}p_{\mu} \int d^{3}q |\psi_{\vec{p}}(\vec{q})|^{2} f_{\alpha}(\vec{p_{\alpha}}, x) f_{\beta}(\vec{p_{\beta}}, x)$$

$$E\frac{dN_{B}(\vec{p})}{d^{3}p} = \int d\sigma^{\mu}p_{\mu} \int d^{3}q_{1}d^{3}q_{2} |\psi_{\vec{p}}(\vec{q_{1}}, \vec{q_{2}})|^{2} f_{\alpha}(\vec{p_{\alpha}}, x) f_{\beta}(\vec{p_{\beta}}, x) f_{\gamma}(\vec{p_{\gamma}}, x)$$
Quark phase space density

Already assumed: rare process, small binding energy, factorization of 2-parton distribution functions, slowly-varying quark spatial-distributions, same hypersurface

Assume quark momentum distribution:

$$f_i(\vec{p_\perp}) \equiv rac{dN}{d^3p} = g_i(p_\perp) \left[1 + 2v_{2,i}(p_\perp)\cos(2\phi)
ight]$$

Yield simple relations of quark and hadron flow:

$$v_{2,M}(p_{\perp}) \approx 2v_{2,q}(\frac{p_{\perp}}{2}) \qquad v_{2,B}(p_{\perp}) \approx 3v_{2,q}(\frac{p_{\perp}}{3})$$



## v<sub>2</sub> and quark number scaling



- Mass ordering at low  $p_T$  driven by the hydrodynamic pressure gradient.
- Baryon meson ordering in the high  $p_T$  region.
- NCQ scaling observed in a wide range of KE<sub>T</sub> indicates the dominance of partonic degrees of freedom.

#### Elliptic flow in default mode



- Mass ordering exists in the low p<sub>T</sub> region
- No baryon meson grouping of v<sub>2</sub>
- NCQ scaling doesn't exist

### Violations to NCQ scaling



- Possible reasons:
  - Narrow wave function width Phys. Lett. B 618:77 (2005)
  - Resonance decay Phys. Rev. C 71:041901 (2005)
  - Higher Fock states contribution J. Phys. G 32:S135 (2006)
  - Space-Momentum correlations at freeze-out
  - Phase-space density Phys. Rev. C 93, 034908 (2016)

*Nucl. Phys. A* 749:268 (2005) , nucl-th/0408044, nuclth/0505061, *Phys. Rev. C* 68:034904 (2003) , *Phys. Rev. C* 70:024901 (2004)

#### Quark hadron flow relations



- Constituent quark flow very close to the formed hadron, different from the amplification behavior expected in coalescence model.
- Hadron flow can not be reverted to the quark flow through simple NCQ rule while NCQ scaling exists!